

Mass Balance

Mass balance for flow system

- For flow problem, apply mass conservation over all the system or part of it.

$$\text{input} = \text{output} + \text{accumulation}$$

- In case of flow problems, we usually deal with flow rates and steady state condition,

$$\text{rate of input} = \text{rate of output (steady state)}$$

Look to the following flow system

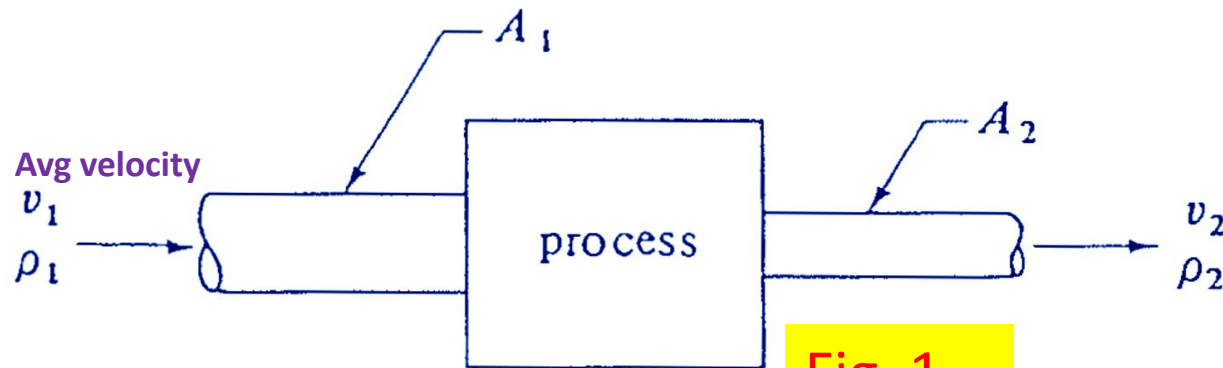


Fig. 1

Mass Balance:

$$m = \rho_1 A_1 v_1 = \rho_2 A_2 v_2 \quad \text{Check units}$$

Note:

Mass velocity or mass flux G

$$G = \rho v \quad \text{with units} \quad \text{kg/s m}^2$$

Control Volume

Control volume for flow through a conduit.

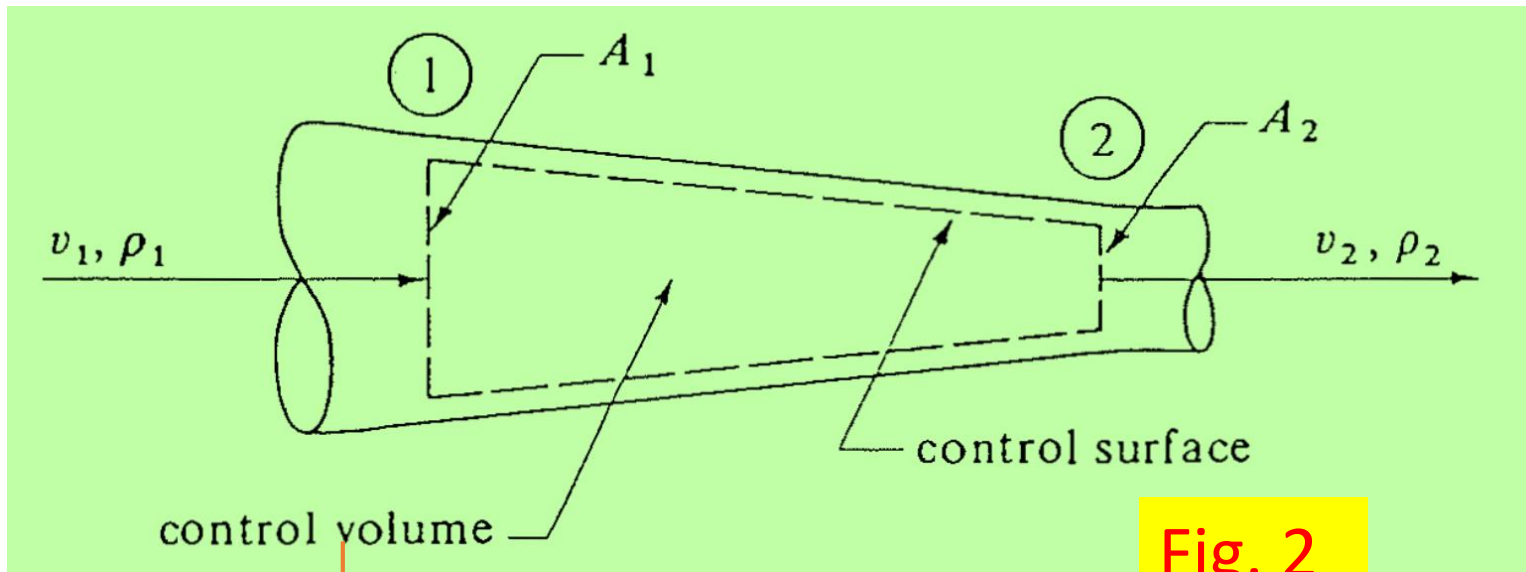


Fig. 2

a region fixed in space through which the fluid flows.

Note:

In most problems the control surface is taken as the wall of the conduit.

Mass balance Eq. over a control volume

In deriving the general equation for the overall balance of the property mass, the law of conservation of mass may be stated as follows for a control volume where no mass is being generated.

$$\begin{aligned} &\left(\begin{array}{l} \text{rate of mass output} \\ \text{from control volume} \end{array} \right) - \left(\begin{array}{l} \text{rate of mass input} \\ \text{from control volume} \end{array} \right) \\ &\quad + \left(\begin{array}{l} \text{rate of mass accumulation} \\ \text{in control volume} \end{array} \right) = 0 \quad (\text{rate of mass generation}) \\ &\hspace{20em} \dots\dots\dots (1) \end{aligned}$$

Assume a general control volume as shown below, and focus on a differential element dA .

Flow through a differential area dA on a control surface.

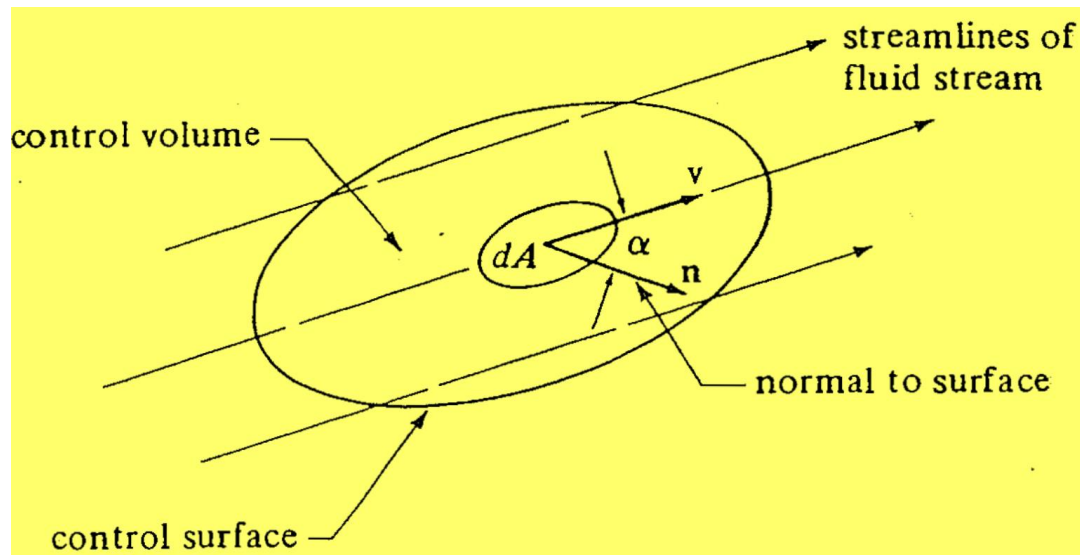


Fig. 3

the rate of mass efflux from this element $= (\rho v)(dA \cos \alpha)$, where $(dA \cos \alpha)$ is the area dA projected in a direction normal to the velocity vector \mathbf{v} , α is the angle between the velocity vector \mathbf{v} and the outward-directed unit normal vector \mathbf{n} to dA

From vector algebra we recognize that $(\rho \mathbf{v})(dA \cos \alpha)$ is the scalar or dot product $\rho(\mathbf{v} \cdot \mathbf{n}) dA$. If we now integrate this quantity over the entire control surface A we have the net outflow of mass across the control surface, or the net mass efflux in kg/s from the entire control volume V :

$\rho \mathbf{v}$ = mass

velocity \mathbf{v} 'flux'

Scalar or dot product

$(\mathbf{v} \cdot \mathbf{n}) = \|\mathbf{v}\| \|\mathbf{n}\| \cos \alpha$

$= v \cdot 1 \cos \alpha = v \cos \alpha$

$$\left(\begin{array}{l} \text{net mass efflux} \\ \text{from control volume} \end{array} \right) = \iint_A v \rho \cos \alpha dA = \iint_A \rho(\mathbf{v} \cdot \mathbf{n}) dA$$

.....(2)

Rate of accumulation

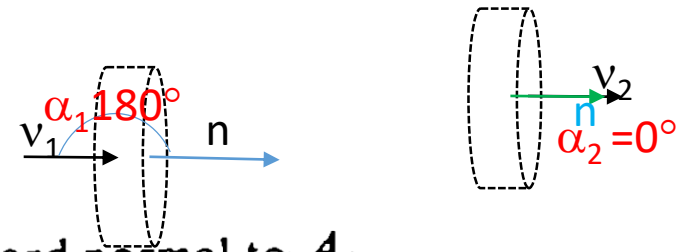
$$\left(\begin{array}{l} \text{rate of mass accumulation} \\ \text{in control volume} \end{array} \right) = \frac{\partial}{\partial t} \iiint_V \rho dV = \frac{dM}{dt}$$

.....(3)

Substituting Eq.^s (2) and (3) into (1) leads to general form of overall mass balance.

$$\iint_A \rho(\mathbf{v} \cdot \mathbf{n}) dA + \frac{\partial}{\partial t} \iiint_V \rho dV = 0 \quad \dots\dots\dots(4)$$

- Apply this eq. to Fig. 2



where all the flow inward is normal to A_1 and outward normal to A_2

When the velocity v_2 leaving is normal to A_2 , the

angle α_2 between the normal to the control surface and the direction of the velocity is 0°

and $\cos \alpha_2 = 1.0$. Where v_1 is directed inward, $\alpha_1 > \pi/2$ α_1 is 180° ($\cos \alpha_1 = -1.0$)

Since α_2 is 0° and α_1 is 180° , using Eq.(4)

$$\begin{aligned} \iint_A v\rho \cos \alpha dA &= \iint_{A_2} v\rho \cos \alpha_2 dA + \iint_{A_1} v\rho \cos \alpha_1 dA \\ &= v_2 \rho_2 A_2 - v_1 \rho_1 A_1 \end{aligned} \quad \dots\dots\dots(5)$$

For steady state, $dM/dt = 0$ in Eq. (3) And (4) becomes

$$m = \rho_1 v_1 A_1 = \rho_2 v_2 A_2$$

Notes: material balance can be done over species or components

$$\underbrace{m_{i2}}_{\text{out}} - \underbrace{m_{i1}}_{\text{in}} + \underbrace{\frac{dM_i}{dt}}_{\text{accumulation}} = \underbrace{R_i}_{\text{generation}} \dots\dots\dots(5a)$$

i means component i in multicomponent system

Average Velocity to Use in Overall Mass Balance

If the velocity is not constant but varies across the surface area, an average or bulk velocity is defined by

$$v_{av} = \frac{1}{A} \iint_A v \, dA$$

for a surface over which v is normal to A and the density ρ is assumed constant.

Example 1

For the case of incompressible flow (ρ is constant) through a circular pipe of radius R , the velocity profile is parabolic for laminar flow as follows:

$$v = v_{\max} \left[1 - \left(\frac{r}{R} \right)^2 \right]$$

where v_{\max} is the maximum velocity at the center where $r = 0$ and v is the velocity at a radial distance r from the center. Derive an expression for the average or bulk velocity v_{av} to use in the overall mass-balance equation.

Solution

The average velocity is represented by

$$v_{av} = \frac{1}{A} \iint_A v \, dA$$

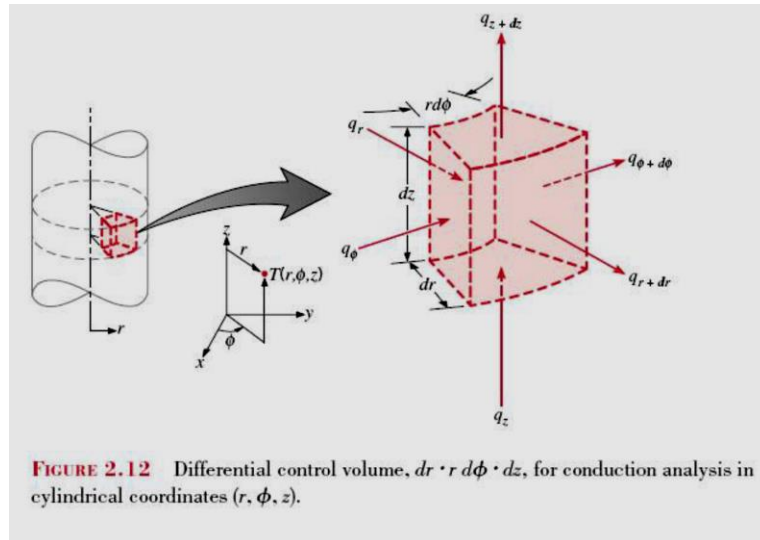
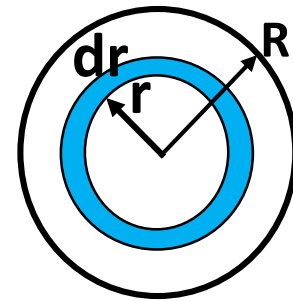
In Cartesian coordinates dA is $dx \, dy$.

However, using polar coordinates which are more appropriate for a pipe, $dA = r \, dr \, d\theta$,

where θ is the angle in polar coordinates. Substituting $dA = r \, dr \, d\theta$, and $A = \pi R^2$ and v in above equation and integrating

$$\begin{aligned} v_{av} &= \frac{1}{\pi R^2} \int_0^{2\pi} \int_0^R v_{max} \left[1 - \left(\frac{r}{R} \right)^2 \right] r \, dr \, d\theta \\ &= \frac{v_{max}}{\pi R^4} \int_0^{2\pi} \int_0^R (R^2 - r^2) r \, dr \, d\theta \\ &= \frac{v_{max}}{\pi R^4} (2\pi - 0) \left(\frac{R^4}{2} - \frac{R^4}{4} \right) \end{aligned}$$

$$v_{av} = \frac{v_{max}}{2}$$



OVERALL ENERGY BALANCE

- Apply the principle of energy conservation to the control volume in the same manner as mass conservation. Begin with 1st law of Thermodynamics

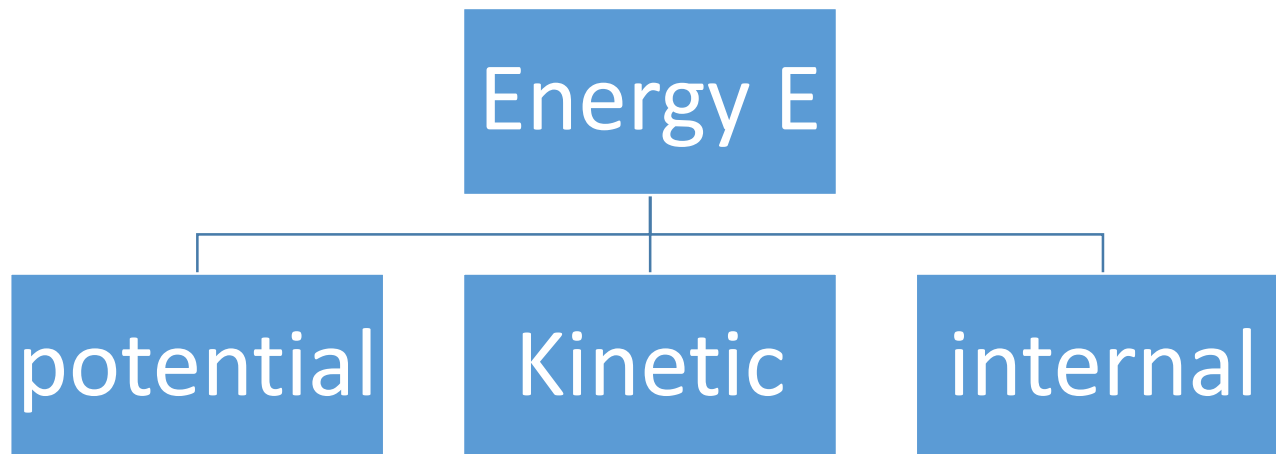
$$\Delta E = Q - W$$

where E is the total energy per unit mass of fluid,
 Q is the heat *absorbed* per unit mass of fluid,
and W is the work of all kinds done per unit mass of fluid
upon the surroundings.

Derivation of Overall Energy-Balance Equation

$$\begin{aligned} &\text{rate of entity output} - \text{rate of entity input} \\ &\quad + \text{rate of entity accumulation} = 0 \\ &\quad \dots\dots\dots(6) \end{aligned}$$

The energy E present within a system can be classified in three ways.



The total energy of the fluid per unit mass is then

$$E = U + \frac{v^2}{2} + zg \quad (\text{SI})$$

$$\left(\begin{array}{c} \text{rate of energy accumulation} \\ \text{in control volume} \end{array} \right) = \frac{\partial}{\partial t} \iiint_V \left(U + \frac{v^2}{2} + zg \right) \rho \, dV \quad \dots\dots\dots(7)$$

Notes:

- # The mass added or removed from the system carries internal, kinetic, and potential energy.
- # In addition, energy is transferred when mass flows into and out of the control volume.
- # pressure–volume work per unit mass fluid is pV .
- # $H = U + pV$
- # the total energy carried with a unit mass is $(H + v^2/2 + zg)$
- # For a small area dA on the control surface the rate of energy efflux is $(H + v^2/2 + zg)(\rho v)(dA \cos \alpha)$

$$\left(\begin{array}{c} \text{net energy efflux} \\ \text{from control volume} \end{array} \right) = \iint_A \left(H + \frac{v^2}{2} + zg \right) (\rho v) \cos \alpha \, dA \dots\dots\dots(8)$$

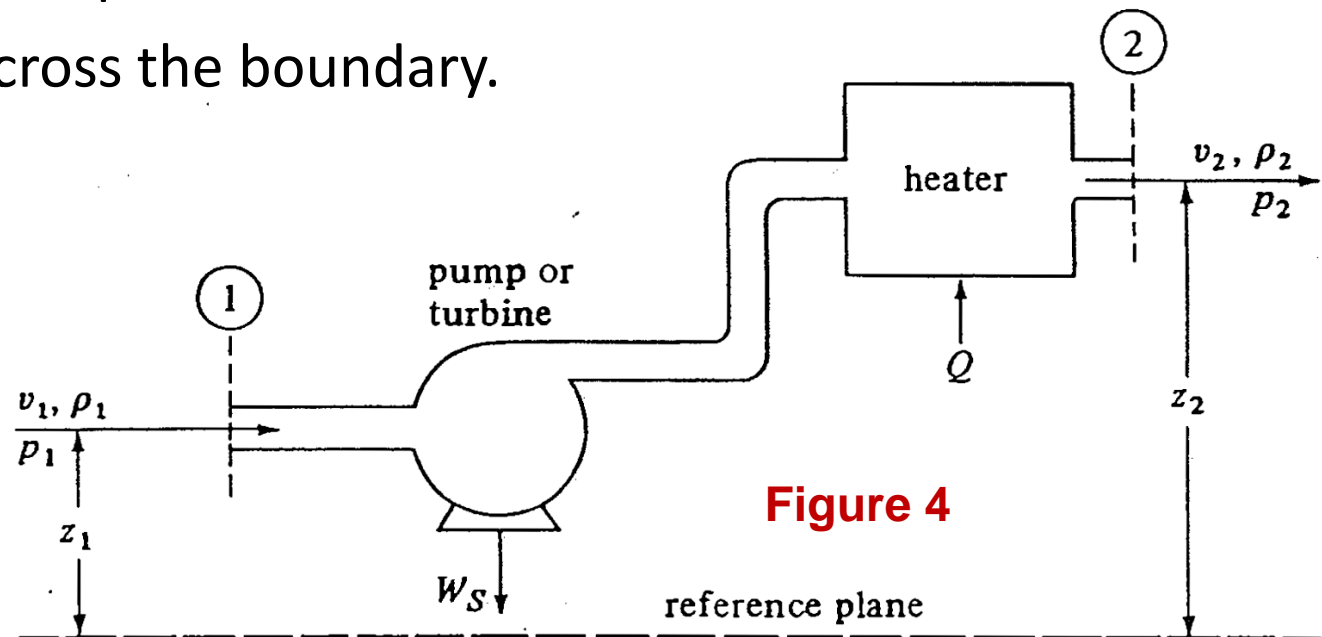
To obtain the overall energy balance, we substitute Eqs. (8) And (7) into (6) and equate the resulting equation to $q - \dot{W}_S$

$$\iint_A \left(H + \frac{v^2}{2} + zg \right) (\rho v) \cos \alpha \, dA + \frac{\partial}{\partial t} \iiint_V \left(U + \frac{v^2}{2} + zg \right) \rho \, dV = q - \dot{W}_S \dots\dots\dots(9)$$

Overall Energy Balance for Steady-State Flow System

Consider the following **figure 4** with the following assumptions:

- Steady-state system.
- Single inlet and outlet.
- Ignore inlet and exit height z , density ρ and enthalpy H variations. Eq. 9 becomes
- 1-d flow cross the boundary.



To determine the overall energy balance for the common system shown in **figure 4**, we can consider the following: since the angle between the velocity vector and unit normal vector is $\alpha = 0$ and since the accumulation term is 0 at steady state, **Eq. 9** can be reduced to

$$\iint_A \left(H + \frac{v^2}{2} + zg \right) (\rho v) \cos \alpha \, dA + \frac{\partial}{\partial t} \iiint_V \left(U + \frac{v^2}{2} + zg \right) \rho \, dV = q - \dot{W}_S$$

.....(9)

$$H_2 m_2 - H_1 m_1 + \frac{m_2 (v_2^3)_{av}}{2v_{2\ av}} - \frac{m_1 (v_1^3)_{av}}{2v_{1\ av}} + gm_2 z_2 - gm_1 z_1 = q - \dot{W}_s \dots\dots\dots(9)$$

For steady state, $m_1 = \rho_1 v_{1\ av} A_1 = m_2 = m$.

Dividing through by m so that the equation is on a unit mass basis,

$$H_2 - H_1 + \frac{1}{2} \left[\frac{(v_2^3)_{av}}{v_{2\ av}} - \frac{(v_1^3)_{av}}{v_{1\ av}} \right] + g(z_2 - z_1) = Q - W_s \dots\dots\dots(10)$$

The term $(v^3)_{av}/(2v_{av})$ can be replaced by $v_{av}^2/2\alpha$.

where α is the kinetic-energy velocity correction factor and is equal to $v_{av}^3/(v^3)_{av}$

Eq. 10 can be rewritten as

$$H_2 - H_1 + \frac{1}{2\alpha} (v_{2\ av}^2 - v_{1\ av}^2) + g(z_2 - z_1) = Q - W_s \dots\dots\dots(10)$$

Note: Kinetic energy term

$$\text{kinetic energy} = \iint_A \left(\frac{v^2}{2} \right) (\rho v) \cos \alpha \, dA \dots\dots\dots(11)$$

Assume $\rho = \text{cons.}$ And set $\cos \alpha = 1$.

Then multiplying the numerator and denominator by $v_{av} A$ noting that $m = \rho v_{av} A$, Eq.(11) becomes

$$\frac{\rho}{2} \iint_A (v^3) \, dA = \frac{\rho v_{av} A}{2 v_{av} A} \iint_A (v^3) \, dA = \frac{m}{2 v_{av} A} \iint_A (v^3) \, dA \dots\dots\dots(12)$$

Dividing through by m so that Eq.(12) becomes per unit mass

$$\left(\frac{1}{2 v_{av}} \right) \frac{1}{A} \iint_A (v^3) \, dA = \frac{(v^3)_{av}}{2 v_{av}} = \frac{v_{av}^2}{2 \alpha} \dots\dots\dots(13)$$

where α is defined as $\alpha = \frac{v_{av}^3}{(v^3)_{av}}$

and $(v^3)_{av}$ is defined as follows:

$$(v^3)_{av} = \frac{1}{A} \iint_A (v^3) dA \dots\dots\dots(14)$$

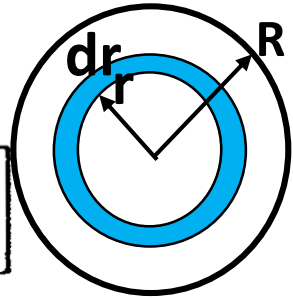
The local velocity v varies across the cross-sectional area of a pipe. To evaluate $(v^3)_{av}$ and, hence, the value of α , we must have an equation relating v as a function of position in the cross-sectional area.

For laminar regime

- Combine $v_{av} = \frac{v_{max}}{2}$ and $v = v_{max} \left[1 - \left(\frac{r}{R} \right)^2 \right]$
 $v = 2v_{av} \left[1 - \left(\frac{r}{R} \right)^2 \right]$
- Substitute in eq. 14 and noting that $A = \pi R^2$ and $dA = r dr d\theta$

$$\begin{aligned} (v^3)_{av} &= \frac{1}{\pi R^2} \int_0^{2\pi} \int_0^R \left[2v_{av} \left(1 - \frac{r^2}{R^2} \right) \right]^3 r dr d\theta \\ &= \frac{(2\pi)2^3 v_{av}^3}{\pi R^2} \int_0^R \frac{(R^2 - r^2)^3}{R^6} r dr = \frac{16v_{av}^3}{R^8} \int_0^R (R^2 - r^2)^3 r dr \end{aligned}$$

- On integration and rearrangement*



$$\begin{aligned}
 (v^3)_{av} &= \frac{16v_{av}^3}{R^8} \int_0^R (R^6 - 3r^2R^4 + 3r^4R^2 - r^6)r \, dr \\
 &= \frac{16v_{av}^3}{R^8} \left(\frac{R^8}{2} - \frac{3}{4} R^8 + \frac{1}{2} R^8 - \frac{1}{8} R^8 \right) \\
 &= 2v_{av}^3
 \end{aligned}$$

$\therefore \alpha$ becomes

$$\alpha = \frac{v_{av}^3}{(v^3)_{av}} = \frac{v_{av}^3}{2v_{av}^3} = 0.50$$

Hence, for laminar flow the value of α to use in the kinetic-energy term of Eq. (10) is 0.5

Note: for turbulent flow α is taken 1 for details see the text.

Example

Initially, a tank contains 500 kg of salt solution containing 10% salt. At point (1) in the control volume in Fig. 4, a stream enters at a constant flow rate of 10 kg/h containing 20% salt. A stream leaves at point (2) at a constant rate of 5 kg/h. The tank is well stirred. Derive an equation relating the weight fraction w_A of the salt in the tank at any time t in hours.

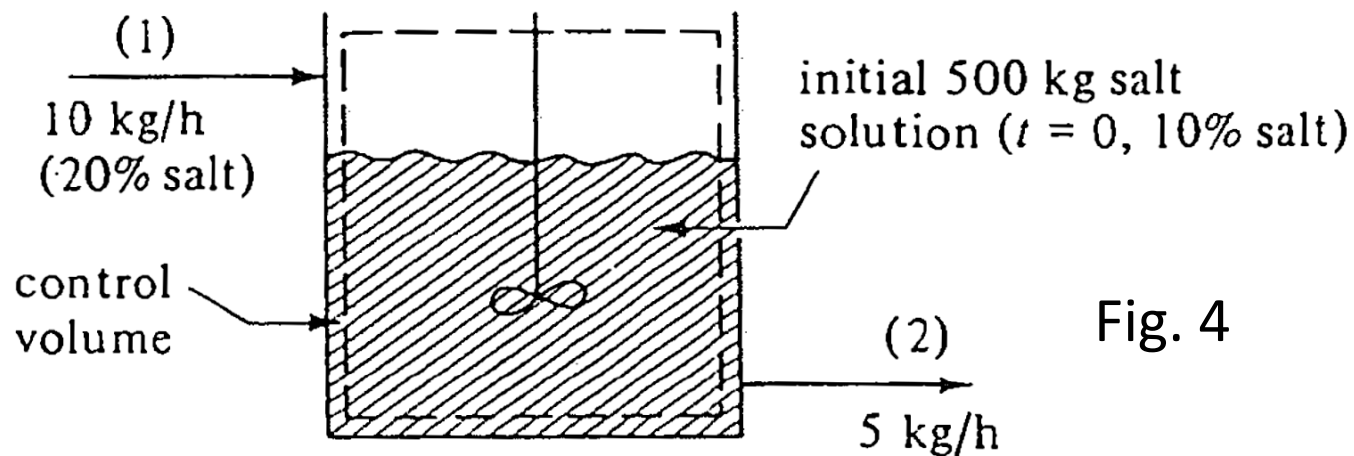


Fig. 4

Solution

$$\iint_A v \rho \cos \alpha \, dA = m_2 - m_1 = 5 - 10 = -5 \text{ kg solution/h}$$

$$\frac{\partial}{\partial t} \iiint_V \rho \, dV = \frac{dM}{dt}$$

- Hence the total mass bal. eq. becomes:

$$-5 + \frac{dM}{dt} = 0$$

$$\int_{M=500}^M dM = 5 \int_{t=0}^t dt$$

$$M = 5t + 500 \quad \text{.....(a)}$$

- Now, make component bal., bal over salt

$$\iint_A v \rho \cos \alpha dA = (5)w_A - 10(0.20) = 5w_A - 2 \text{ kg salt/h}$$

$$\frac{\partial}{\partial t} \iiint_V \rho dV = \frac{d}{dt} (Mw_A) = \frac{M}{dt} \frac{dw_A}{dt} + w_A \frac{dM}{dt} \text{ kg salt/h}$$

- Hence salt bal. eq. becomes {component Bal. Eq.}

$$5w_A - 2 + M \frac{dw_A}{dt} + w_A \frac{dM}{dt} = 0 \quad \dots\dots\dots(b)$$

- Substitute M from eq.(a) in eq. (b)

$$5w_A - 2 + (500 + 5t) \frac{dw_A}{dt} + w_A \frac{d(500 + 5t)}{dt} = 0$$

- On rearrangement and integration

$$5w_A - 2 + (500 + 5t) \frac{dw_A}{dt} + 5w_A = 0$$

$$\int_{w_A=0.10}^{w_A} \frac{dw_A}{2 - 10w_A} = \int_{t=0}^t \frac{dt}{500 + 5t}$$

$$-\frac{1}{10} \ln \left(\frac{2 - 10w_A}{1} \right) = \frac{1}{5} \ln \left(\frac{500 + 5t}{500} \right)$$

$$w_A = -0.1 \left(\frac{100}{100 + t} \right)^2 + 0.20$$

Comments

- ✓ Overall or macroscopic material (or energy or momentum) gives an idea about the system from outside the enclosure.
- ✓ Therefore, overall balances (material; energy; momentum) do not tell us the details of what happens inside the system.
- ✓ The previous discussion is useful for the next step “momentum balance and shell momentum balance’.
- ✓ The shell momentum balance will be made in order to obtain the details of what happens inside the system.

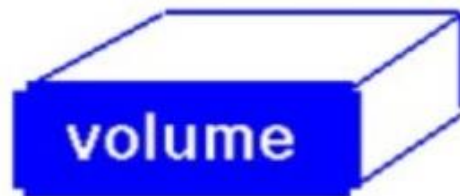
Scalars and Vectors

A scalar quantity has only **magnitude**.

A vector quantity has both **magnitude** and **direction**.

Scalar Quantities

length, area, volume
speed
mass, density
pressure
temperature
energy, entropy
work, power



Vector Quantities

displacement
velocity
acceleration
momentum
force
lift, drag, thrust
weight

