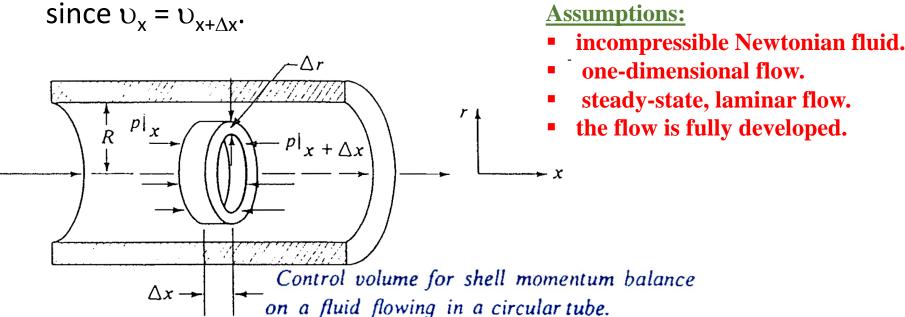
# Shell Momentum Balance

INSIDE PIPES & FALLING FILM

## SHELL MOMENTUM BALANCE AND VELOCITY PROFILE IN LAMINAR FLOW

- Shell Momentum Balance Inside a Pipe
- Consider a flow of fluid inside a circular conduit or pipe.
- Assume a control volume as shown in the figure.
- Assume <u>fully developed</u> flow i.e velocity,  $\upsilon = f(r) \neq f(x)$ . No net momentum flux across the annular volume at x and  $x + \Delta x$



At steady state the conservation of momentum, Eq.(3) previous slides. becomes as follows:

The pressure forces become

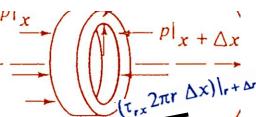
pressure forces = 
$$pA \mid_x - pA \mid_{x+\Delta x} = p(2\pi r \Delta r) \mid_x - p(2\pi r \Delta r) \mid_{x+\Delta x}$$
 (2)

#### The shear force or drag force

The shear force or drag force acting on the cylindrical surface at the radius r is the shear stress  $\tau_{rx}$  times the area  $2\pi r \Delta x$ . This can also be considered as the rate of momentum flow into the cylindrical surface of the shell as described by

$$\tau_{yz} = \frac{kg \cdot m/s}{m^2 \cdot s} = \frac{momentum}{m^2 \cdot s}$$

#### ... The net rate of momentum efflux is



the rate of momentum out - rate of momentum in

Or

net efflux = 
$$(\tau_{rx} 2\pi r \Delta x)|_{r+\Delta r} - (\tau_{rx} 2\pi r \Delta x)|_{r}$$
 (3)

#### Note:

The net momentum flux across the annular volume at x and  $x + \Delta x = 0$ . Now, substituting eqs. (2) and (3) into eq. (1) and making rearrangement Hence,

$$\frac{(r\tau_{rx})|_{r+\Delta r} - (r\tau_{rx})|_{r}}{\Delta r} = \frac{r(p|_{x} - p|_{x+\Delta x})}{\Delta x} \tag{4}$$

In fully developed flow, the pressure gradient  $(\Delta p/\Delta x)$  is constant and becomes  $(\Delta p/L)$ , where  $\Delta p$  is the pressure drop for a pipe of length L. Letting  $\Delta r$  approach zero, we obtain

$$\frac{d(r\tau_{rx})}{dr} = \left(\frac{\Delta p}{L}\right)r \qquad (5)$$

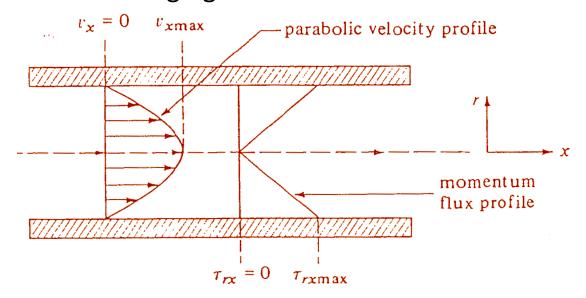
Separating variables and integrating,

$$\tau_{rx} = \left(\frac{\Delta p}{L}\right) \frac{r}{2} + \frac{C_1}{r} \qquad (6)$$

The constant of integration  $C_1$  must be zero if the momentum flux is not infinite at r = 0.

$$\tau_{rx} = \left(\frac{\Delta p}{2L}\right) r = \frac{p_0 - p_L}{2L} r \qquad \dots \tag{7}$$

This means that the momentum flux varies linearly with the radius as shown in the following figure.



Substituting Newton's law of viscosity in the previous equation,

$$\tau_{rx} = -\mu \, \frac{dv_x}{dr} \qquad \dots \tag{8}$$

We obtain

$$\frac{dv_x}{dr} = -\frac{p_0 - p_L}{2\mu L} r \tag{9}$$

Integrating using the boundary condition that at the wall,  $v_x = 0$  at r = R, we obtain the equation for the velocity distribution.

$$v_x = \frac{p_0 - p_L}{4\mu L} R^2 \left[ 1 - \left(\frac{r}{R}\right)^2 \right] \qquad (10)$$

This result shows that the velocity distribution is parabolic as indicated in the previous figure.

#### To obtain avg velocity

Use the average concept equation as given earlier.

$$v_{x \text{ av}} = \frac{1}{A} \iint_{A} v_{x} dA = \frac{1}{\pi R^{2}} \int_{0}^{2\pi} \int_{0}^{R} v_{x} r dr d\theta = \frac{1}{\pi R^{2}} \int_{0}^{R} v_{x} 2\pi r dr d\theta$$
(11)

- Note:  $dA = r dr d\theta$  and  $A = \pi R^2$
- Combining eqs. (10) and (11) and integrating

$$v_{x \text{ av}} = \frac{(p_0 - p_L)R^2}{8\mu L} = \frac{(p_0 - p_L)D^2}{32\mu L} \qquad \dots (12)$$

• This eq. is called *Hagen–Poiseuille equation* which relates the pressure drop and the average velocity for laminar flow in a horizontal pipe.

### Max velocity can be found from eq. 10 at r = 0

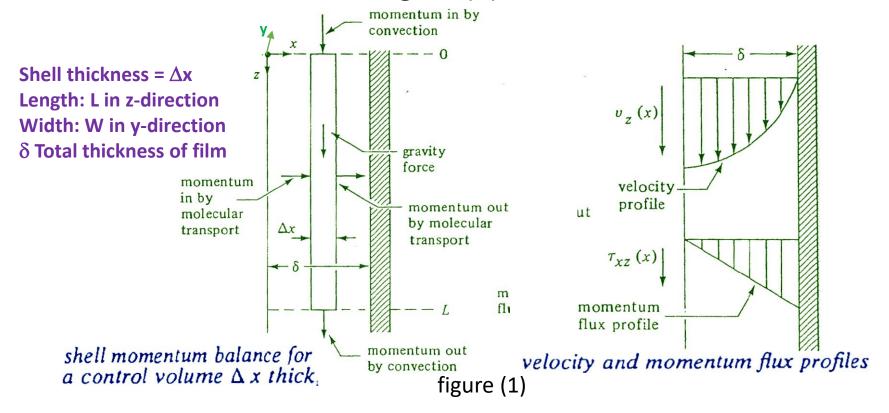
$$v_{x \max} = \frac{p_0 - p_L}{4\mu L} R^2 \dots (13)$$

Combining eqs. (13) and (12), we obtain

$$v_{x \text{ av}} = \frac{v_{x \text{ max}}}{2} \qquad \dots \tag{13}$$

## Shell Momentum Balance for Falling Film

- Meaning: a flow of a fluid as a film in laminar flow down a vertical surface.
- Applications: various phenomena in mass transfer; coatings on surfaces; condensation as filmwise condensation .. etc.
- Control volume: shown in figure (1).



• momentum flux due to molecular transport

net efflux = 
$$LW(\tau_{xz})|_{x+\Delta x} - LW(\tau_{xz})|_{x}$$
 (14)

convective momentum flux

Assume Control volume is far from entrance and exit (no entrance and exit effects on shell flow), hence  $v_{7}(x)$  is independent of z-direction.

net efflux is equal to 0 since  $v_z$  at z = 0 is equal to  $v_z$  at z = L for each value of x.

net efflux = 
$$\Delta x W v_z(\rho v_z)|_{z=L} - \Delta x W v_z(\rho v_z)|_{z=0} = 0$$
 .....(15)

Gravity force acting on the fluid.

gravity force = 
$$\Delta x W L(\rho g)$$
 .....(16)

 Then substitute in momentum conservation eq. at steady state:

$$\Delta x W L(\rho g) = L W(\tau_{xz})|_{x + \Delta x} - L W(\tau_{xz})|_{x} + 0$$
(17)

• Rearranging and letting  $\Delta x \rightarrow 0$ 

$$\frac{\tau_{xz}|_{x+\Delta x} - \tau_{xz}|_{x}}{\Delta x} = \rho g \qquad (18)$$

$$\frac{d}{dx}\tau_{xz} = \rho g \qquad \dots (19)$$

Integrating using the boundary conditions at x = 0,  $\tau_{xz} = 0$ 

at the free liquid surface and at x = x,  $\tau_{xz} = \tau_{xz}$ ,

$$\tau_{xz} = \rho g x \qquad \dots (20)$$

This means the momentum-flux profile is linear as shown in Fig. (1) and the max value at the wall.

Using Newton's Law of viscosity:

$$\tau_{xz} = -\mu \, \frac{dv_z}{dx} \qquad \dots \tag{21}$$

Combining eqs. (20) and (21)

$$\frac{dv_z}{dx} = -\left(\frac{\rho g}{\mu}\right)x \qquad \dots \tag{22}$$

• Separating variables and integrating gives

$$v_z = -\left(\frac{\rho g}{2\mu}\right) x^2 + C_1$$
 (23)

- Using BC. At wall,  $x = \delta$ ;  $v_z = 0$ . hence,  $C_1 = (\rho g/2\mu)\delta^2$
- Then

$$v_z = \frac{\rho g \delta^2}{2\mu} \left[ 1 - \left( \frac{x}{\delta} \right)^2 \right] \qquad (24)$$

This means the velocity profile is parabolic as shown in Fig.(1)

The max velocity occurs at x = 0; hence,

$$v_{z \max} = \frac{\rho g \delta^2}{2u} \qquad (25)$$

Avg velocity as before

$$v_{z \text{ av}} = \frac{1}{A} \iint_{A} v_{z} dA = \frac{1}{W\delta} \int_{0}^{W} \int_{0}^{\delta} v_{z} dx dy = \frac{W}{W\delta} \int_{0}^{\delta} v_{z} dx$$
 (26)

Substituting e. (24) into (26) and integrating

$$v_{z \text{ av}} = \frac{\rho g \delta^2}{3\mu} \qquad \dots (27)$$

Combing eqs. (25) and (27) we obtain,

$$v_{z \text{ av}} = (2/3)v_{z \text{ max}}$$
 (28)

• The volumetric flow rate, q, can be obtained by  $A_c = \delta W$  multiplying the avg velocity times the cross-sectional area

## <u>Note</u>

Often in falling films, the mass rate of flow per unit width of wall  $\Gamma$  in kg/s·m is defined

as  $\Gamma = \rho \delta v_{zav}$  and a Reynolds number is defined as

$$N_{\text{Re}} = \frac{4\Gamma}{\mu} = \frac{4\rho \delta v_{z \text{ av}}}{\mu} \qquad (30)$$

Laminar flow occurs for  $N_{Re}$  < 1200.

Laminar flow with rippling present occurs above a  $N_{Re}$  of 25.

## Example

An oil is flowing down a vertical wall as a film 1.7 mm thick. The oil density is 820 kg/m<sup>3</sup> and the viscosity is 0.20 Pa·s. Calculate the mass flow rate per unit width of wall,  $\Gamma$ , needed and the Reynolds number. Also calculate the average velocity.

## Solution

• Substituting into the definition of  $\Gamma$ 

$$\Gamma = \rho \delta v_{z \, av} = \frac{(\rho \delta) \rho g \delta^2}{3\mu} = \frac{\rho^2 \delta^3 g}{3\mu}$$
$$= \frac{(820)^2 (1.7 \times 10^{-3})^3 (9.806)}{3 \times 0.20} = 0.05399 \text{ kg/s} \cdot \text{m}$$

Check Re

$$N_{\text{Re}} = \frac{4\Gamma}{\mu} = \frac{4(0.05399)}{0.20} = 1.080$$

Hence, the film is in laminar flow.

Avg velocity

$$v_{zav} = \frac{\rho g \delta^2}{3\mu} = \frac{820(9.806)(1.7 \times 10^{-3})^2}{3(0.20)} = 0.03873 \text{ m/s}$$