

# Discussion \_ Heat

Applications on Heat Diffusion  
Equations

1-D, S.S conduction, no internal generation

- Temp distribution + heat transfer rate (planar, cylindrical, spherical)
- thermal resistance

Temp. distribution

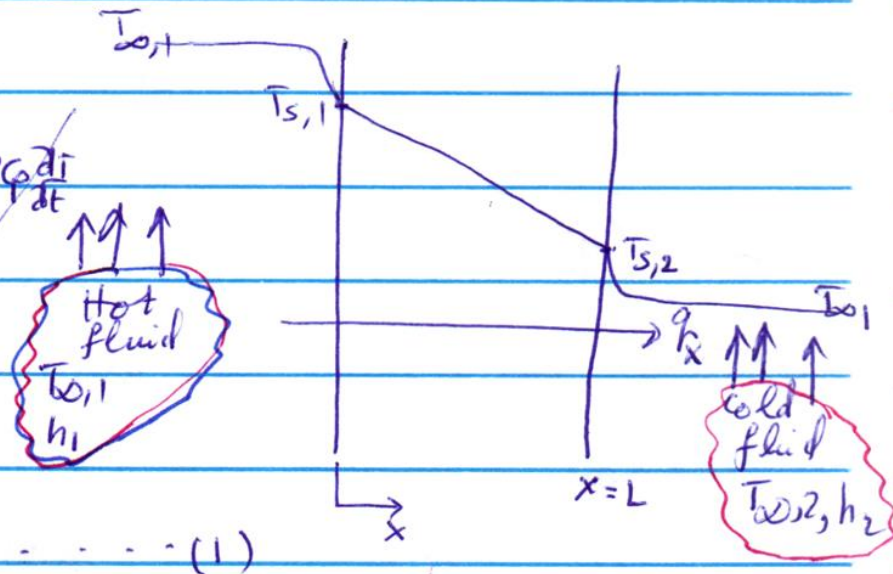
\* The plan wall ✓

$$\frac{\partial}{\partial x} \left( k \frac{\partial T}{\partial x} \right) + \frac{\partial}{\partial y} \left( k \frac{\partial T}{\partial y} \right) + \frac{\partial}{\partial z} \left( k \frac{\partial T}{\partial z} \right) + \dot{q} = \rho c_p \frac{\partial T}{\partial t}$$

Assume  $k = \text{const.}$

$$\frac{d}{dx} \left( k \frac{dT}{dx} \right) = 0$$

Integrating twice



plane wall

$$\frac{d}{dx} \left( k \frac{dT}{dx} \right) = 0$$

..... (1)

plane wall

Integrating twice

$$\frac{dT}{dx} = C_1$$

..... (2)

$$T(x) = C_1 x + C_2$$

..... (3)

To obtain  $C_1, C_2$  we need to apply B.C's

at  $x = 0$

$$T = T_{s,1}$$

~~\*~~

B.C(1)

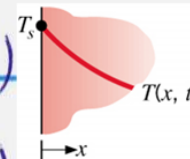
at  $x = L$

$$T(L) = T_{s,2}$$

~~\*~~

B.C(2)

Constant Surface Temperature:



$$T(0, t) = T_s$$

apply BC (1)

$$T_{s,1} = C_2 \quad \therefore C_2 = T_{s,1} \quad \dots (4)$$

apply BC (2)

$$T_{s,2} = C_1 L + T_{s,1} \quad \therefore C_1 = (T_{s,2} - T_{s,1})/L \quad \dots (5)$$

$\therefore$  Temp. distribution

$$T(x) = \frac{(T_{s,2} - T_{s,1})}{L} x + T_{s,1} \quad \text{X} \quad (6)$$

so linear distribution "at the previous conditions".

Note Rearrangement of Eq (6)

$$\frac{T(x) - T_{s,1}}{T_{s,2} - T_{s,1}} = \left( \frac{x}{L} \right) \quad \text{X}$$

To obtain H.T.R

apply Fourier's Law

$$q_x = -k A \frac{dT}{dx} \quad \leftarrow \text{from eq.(2)} \quad \frac{dT}{dx} = C_1 = (T_{s,2} - T_{s,1})/L$$
$$= -k A \frac{T_{s,2} - T_{s,1}}{L}$$

$$q_x'' = \frac{q_x}{A} = -\frac{k}{L} (T_{s,2} - T_{s,1}) = \frac{k}{L} (T_{s,1} - T_{s,2})$$

In case convection B. C.<sup>s</sup>

Follow the same procedure



# Thermal resistance

## Thermal Resistance

$$\therefore q_k = k A \frac{(T_{s,1} - T_{s,2})}{L} = \frac{T_{s,1} - T_{s,2}}{L/kA} = \frac{T_{s,1} - T_{s,2}}{R_{t,cond.}} = \frac{\text{Potential driving}}{\text{Resistance}}$$

$$\therefore R_{t,cond} = L/kA \quad \#$$

Note — Similarly for electrical conduction in the same system.

Ohm's Law provides an electrical resis.

$$R_e = \frac{\overset{\text{electrical potential}}{E_{s,1} - E_{s,2}}}{I} = \frac{L}{\sigma A}$$

electrical conductivity

OR  $1/\Omega \cdot \text{cm}$

$$\sigma = 1/\rho$$

resistivity  $\Omega \cdot \text{cm}$

$$R_e = \frac{\rho L}{A}$$

wire cross section  $\text{cm}^2$

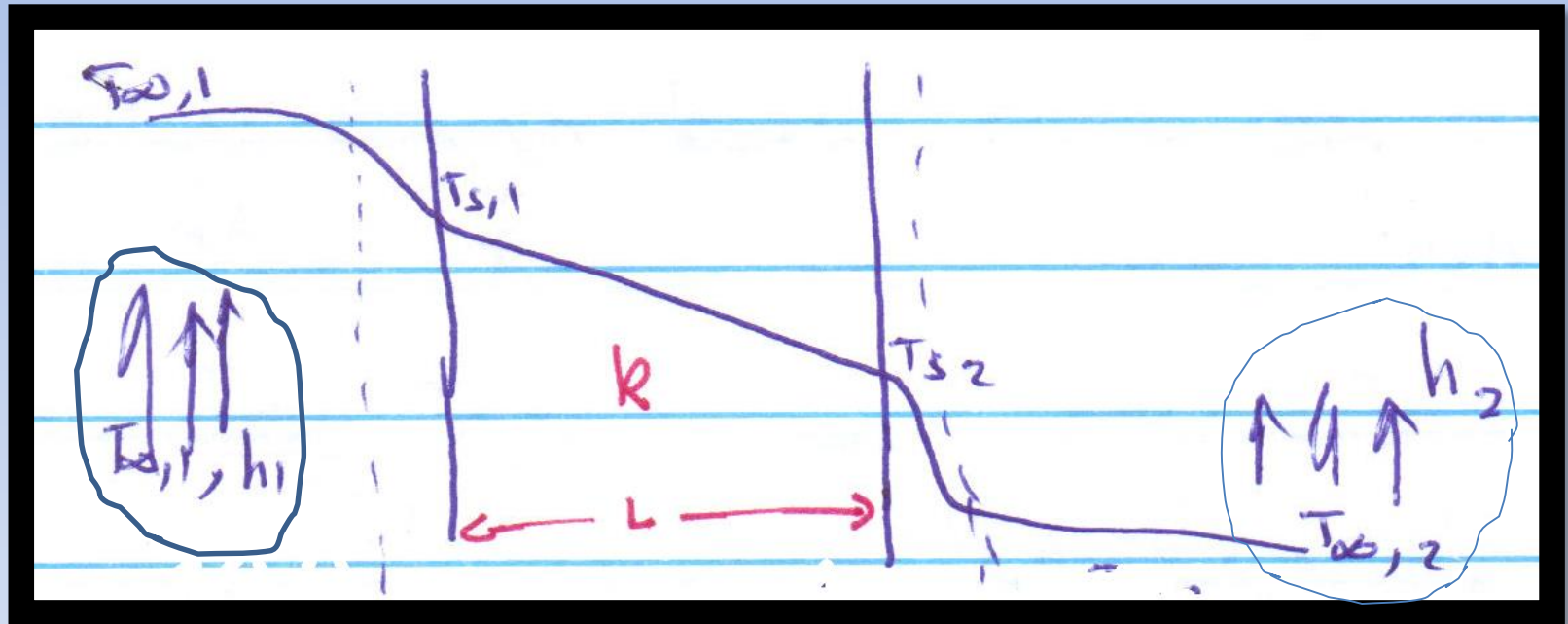
$$\text{Current } I = \frac{E_{s,1} - E_{s,2}}{R_e}$$

For convection Heat Transfer ; Newton's Law

$$q_r = h A (T_s - T_\infty) = \frac{(T_s - T_\infty)}{1/hA} = \frac{(T_s - T_\infty)}{R_{t, \text{conv}}}$$

$$\therefore R_{t, \text{conv}} = 1/hA$$

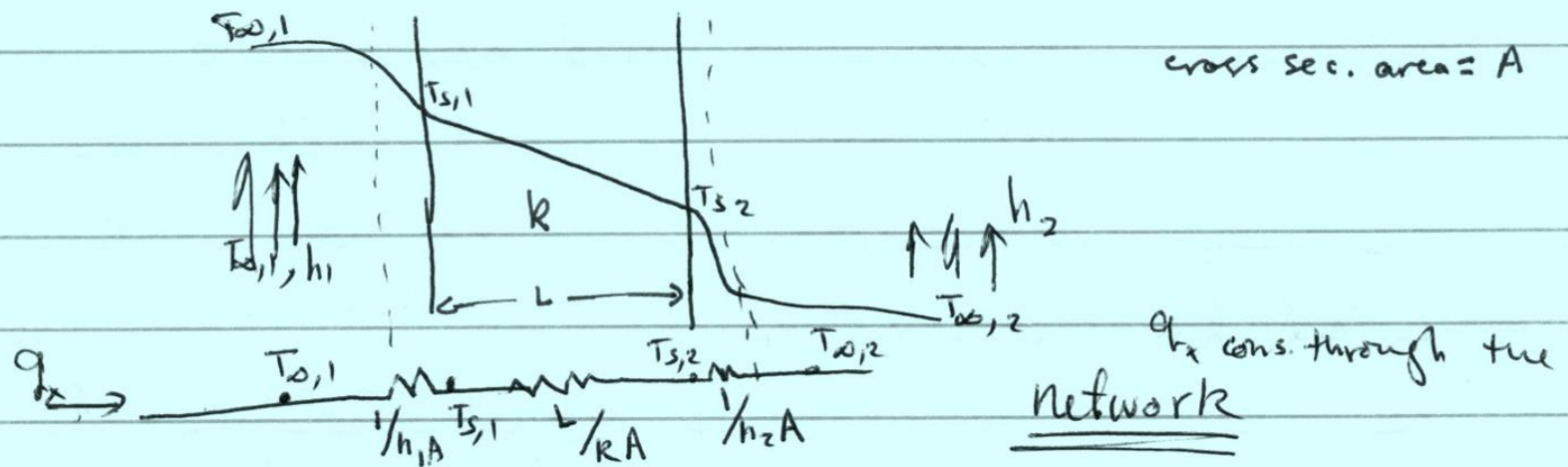
Find the Equivalent thermal circuit  
'Network'





# Equivalent thermal circuit

Now, Circuit representation or equivalent thermal circuit for the plane wall with convection surface conditions as shown below



$$q_x = \frac{T_{\infty,1} - T_{s,1}}{\frac{1}{h_1 A}} = \frac{T_{s,1} - T_{s,2}}{\frac{L}{kA}} = \frac{T_{s,2} - T_{\infty,2}}{\frac{1}{h_2 A}}$$

# Overall temperature difference

In terms the overall temp diff.

$$q_x = \frac{T_{\infty,1} - T_{\infty,2}}{R_{\text{total}}}$$

because the conduction and convection resistances are in series.

$$R_{\text{total}} = \frac{1}{h_1 A} + \frac{L}{kA} + \frac{1}{h_2 A}$$

## Note

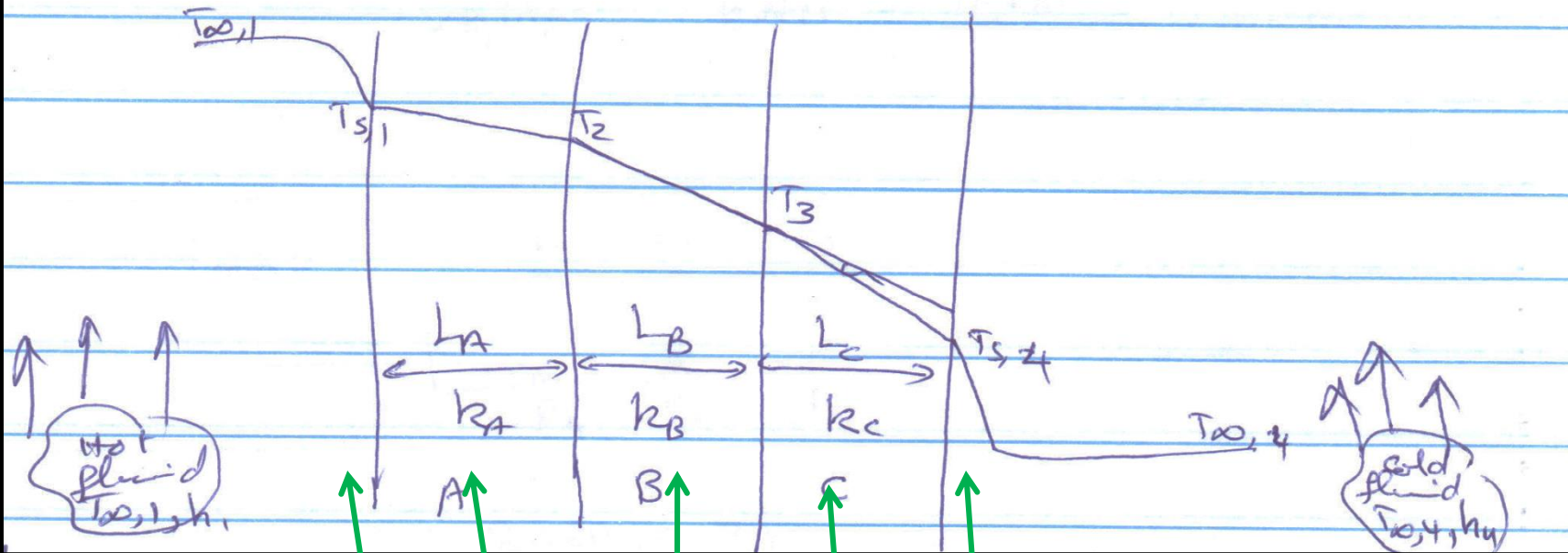
Thermal resis. for rad. could be defined via

$$q_{\text{rad}} = h_r A (T_s - T_{\text{sur}}) = \frac{(T_s - T_{\text{sur}})}{\frac{1}{h_r A}} = \frac{T_s - T_{\text{sur}}}{R_{\text{rad}}}$$

where  $R_{\text{rad}} = \frac{1}{h_r A}$

Beaming in mind that  $R_{\text{rad}}$  and  $R_{\text{conv}}$  act in parallel.

# Composite Wall



Thermal resistance network diagram showing the series combination of resistances:  $\frac{1}{h_1 A}$ ,  $\frac{L_A}{k_A A}$ ,  $\frac{L_B}{k_B A}$ ,  $\frac{L_C}{k_C A}$ , and  $\frac{1}{h_4 A}$ . The temperatures at the nodes are  $T_{\infty,1}$ ,  $T_{s,1}$ ,  $T_2$ ,  $T_3$ ,  $T_{s,4}$ , and  $T_{\infty,4}$ . The heat flux  $q_x$  is indicated entering the first node.

$$q_x = \frac{T_{\infty,1} - T_{\infty,4}}{\sum R} = \frac{T_{\infty,1} - T_{\infty,4}}{\left[ \left( \frac{1}{h_1 A} \right) + \left( \frac{L_A}{k_A A} \right) + \left( \frac{L_B}{k_B A} \right) + \left( \frac{L_C}{k_C A} \right) + \left( \frac{1}{h_4 A} \right) \right]}$$

$$q_x = \frac{T_{\infty,1} - T_{s,1}}{\left( \frac{1}{h_1 A} \right)} = \frac{T_{s,1} - T_2}{\left( \frac{L_A}{k_A A} \right)} = \dots$$

# Overall Heat transfer Coefficient 'U'

## Overall Heat Transfer Coeff. U

\* U can be defined by an expression analogous to Newton's law of cooling

$$Q_x = U A \Delta T \quad \text{--- II}$$

where

$\Delta T$  : overall temp diff.

A : surface area

U overall is related to the total thermal resistance,  $R_t$  which is given in the previous eqs i.e

$$Q_x = \frac{\Delta T}{R_{total}} = \frac{T_{\infty,1} - T_{\infty,2}}{\sum R_t} \quad \text{--- III}$$



Compare the previous eqs

$$\therefore UA = \frac{1}{R_{\text{total}}}$$

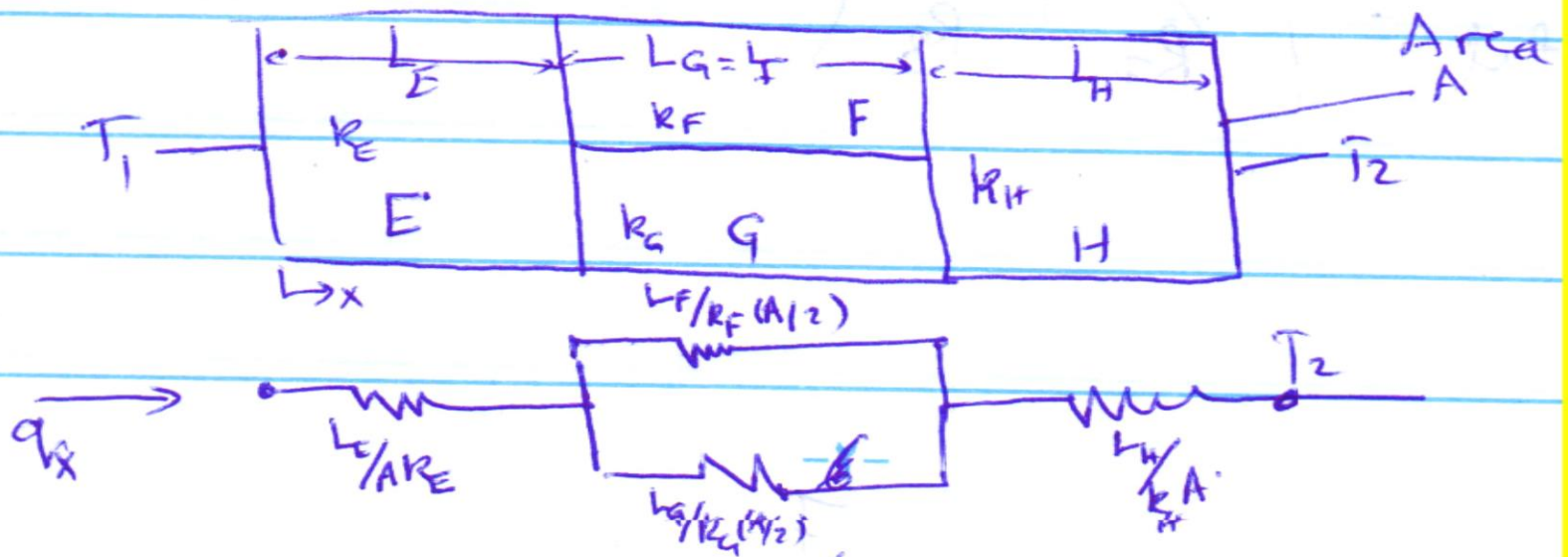
$$\therefore U = \frac{1}{R_{\text{total}} A} = \frac{1}{\left[ \frac{1}{h_1} + \frac{L_A}{k_A} + \frac{L_B}{k_B} + \frac{L_C}{k_C} + \frac{1}{h_2} \right]} \quad \#$$

In general

$$\boxed{R_{\text{tot}} = \sum R_t = \frac{\Delta T}{q} = \frac{1}{UA}} \quad \#$$

# Note

Parallel Configurations  
Series-Parallel configurations





# The Plane Wall

Consider one-dimensional, steady-state conduction in a plane wall of constant  $k$ , with uniform generation, and asymmetric surface conditions:

- Heat diffusion equation (eq. 2.3) :

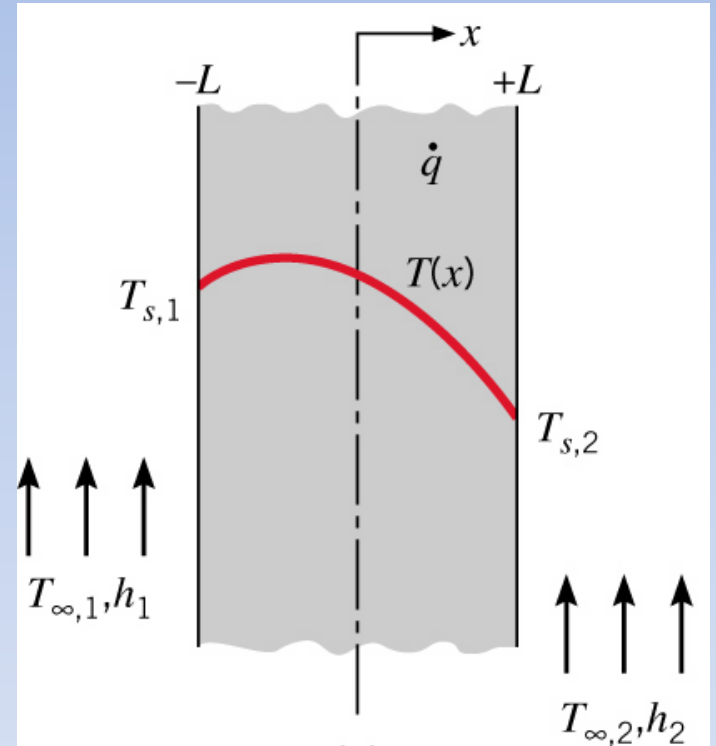
$$\frac{d^2T}{dx^2} + \frac{\dot{q}}{k} = 0$$

- General Solution after taking double integration:

$$T = -\frac{\dot{q}}{2k} x^2 + C_1 x + C_2$$

- Boundary Conditions:

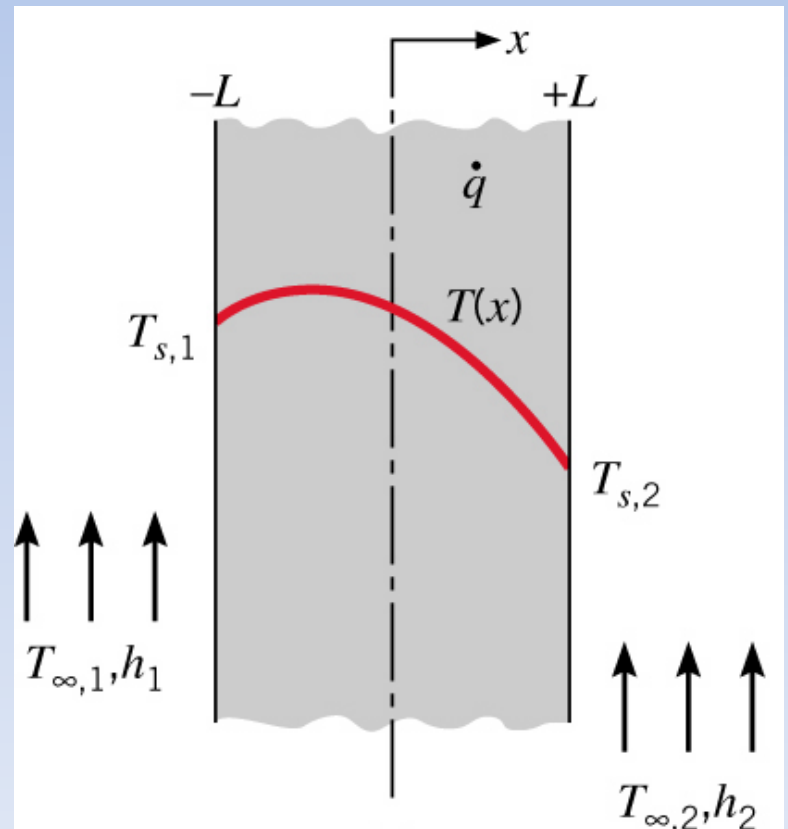
$$T(-L) = T_{s,1}, T(L) = T_{s,2}$$



# Temperature Profile

$$T(x) = \frac{\dot{q}L^2}{2k} \left( 1 - \frac{x^2}{L^2} \right) + \frac{(T_{s,2} - T_{s,1})}{2} \frac{x}{L} + \frac{T_{s,1} + T_{s,2}}{2}$$

- ❖ Profile is parabolic.
- ❖ Heat flux is not independent of  $x$



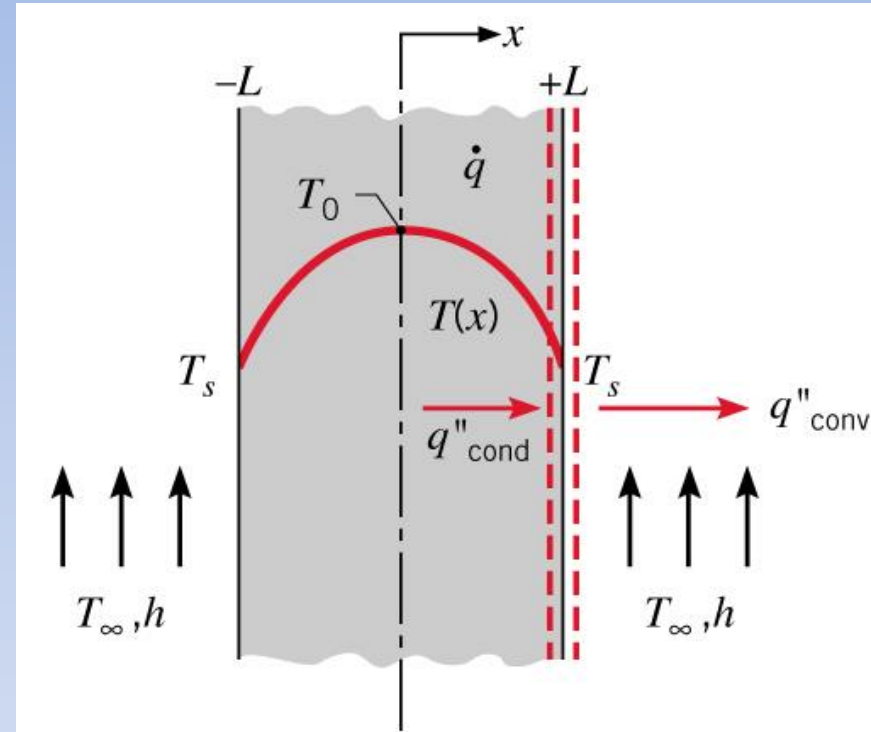
# Symmetrical Distribution

- When both surfaces are maintained at a common temperature,  $T_{s,1} = T_{s,2} = T_s$
- Previous equation becomes

$$T(x) = \frac{\dot{q}L^2}{2k} \left( 1 - \frac{x^2}{L^2} \right) + T_s$$

? What is the location of the maximum temperature?

$$\therefore \frac{T(x) - T_{\max}}{T_s - T_{\max}} = \left( \frac{x}{L} \right)^2$$

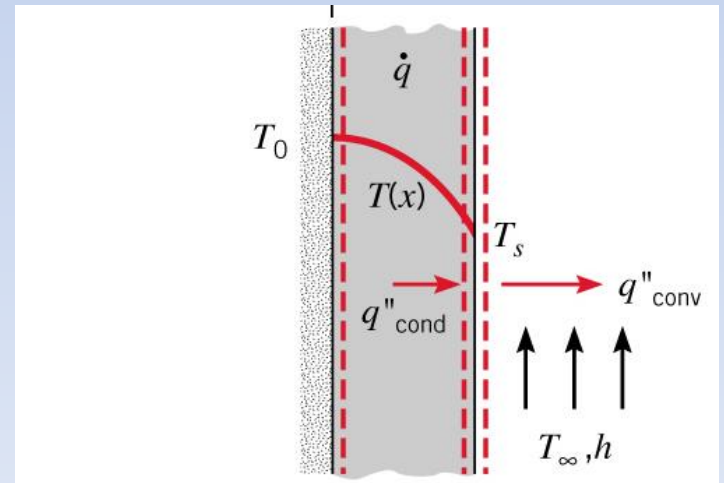
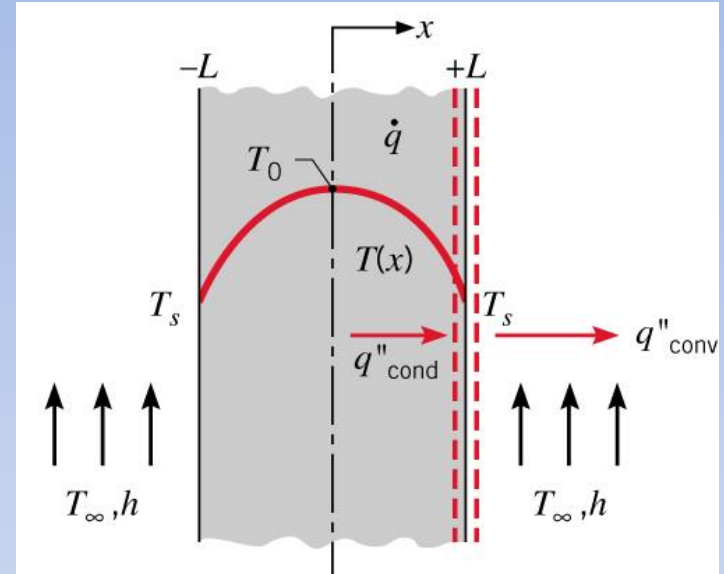


# Symmetrical Distribution

➤ Note that at the plane of symmetry:

$$\left( \frac{dT}{dx} \right)_{x=0} = 0 \Rightarrow q''|_{x=0} = 0$$

❖ Equivalent to adiabatic surface



# Calculation of surface temperature $T_s$

In previous two equations the surface temperature,  $T_s$  is needed.

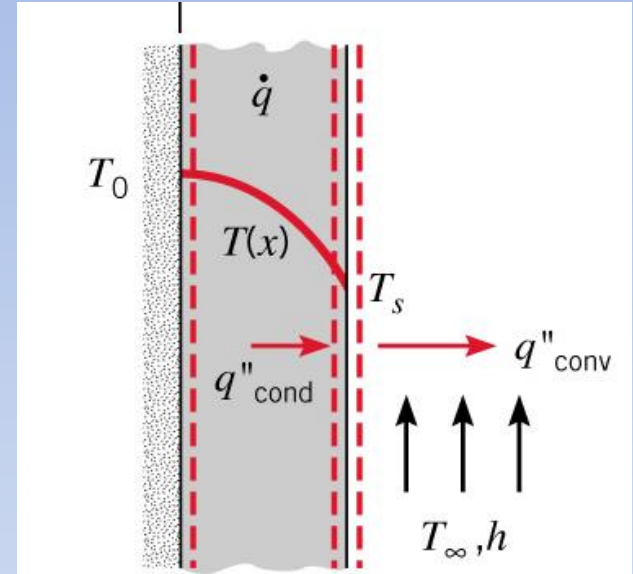
➤ Boundary condition at the wall:

$$-k \left. \frac{dT}{dx} \right|_{x=L} = h(T_s - T_\infty)$$

Substituting  $(dT/dx)_{x=L}$  from equation of temp. distribution i.e

$$T(x) = \frac{\dot{q}L^2}{2k} \left( 1 - \frac{x^2}{L^2} \right) + T_s$$

$$\therefore T_s = T_\infty + \frac{\dot{q}L}{h}$$



**Note: You can obtain the same results by applying energy balance;**

$$\dot{q} AL = hA(T_s - T_\infty)$$

$$T_s = \frac{\dot{q}L}{h} + T_\infty$$

# Example 1

A plane wall of thickness 0.1 m and thermal conductivity 25 W/m·K having uniform volumetric heat generation of 0.3 MW/m<sup>3</sup> is insulated on one side, while the other side is exposed to a fluid at 92°C. The convection heat transfer coefficient between the wall and the fluid is 500 W/m<sup>2</sup>·K. Determine the maximum temperature in the wall.

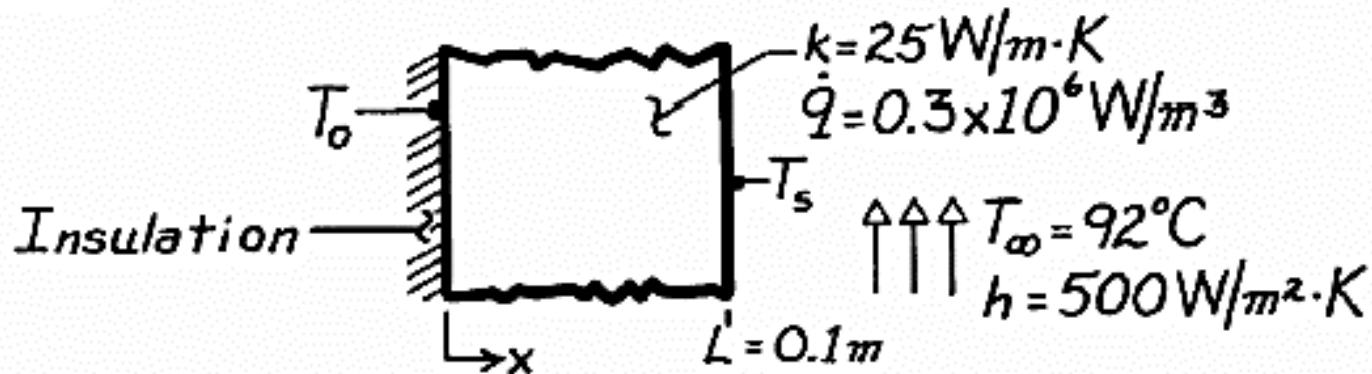


# Solution

**KNOWN:** Plane wall with internal heat generation which is insulated at the inner surface and subjected to a convection process at the outer surface.

**FIND:** Maximum temperature in the wall.

**SCHEMATIC:**



**ASSUMPTIONS:** (1) Steady-state conditions, (2) One-dimensional conduction with uniform volumetric heat generation, (3) Inner surface is adiabatic.

**ANALYSIS:** From previous equations for symmetrical wall, the max temp. is:

$$T_o = \dot{q}L^2 / 2k + T_s. \quad \text{But } T_s \text{ not given}$$

# Solution Continue

The outer surface temperature follows from Previous equation of symmetrical

$$T_s = T_{\infty} + \dot{q}L/h$$

Apply Overall heat balance

$$\dot{E}_{in} - \dot{E}_{out} + \dot{E}_{gen} = \dot{E}_{st} \quad \therefore \dot{E}_{gen} = \dot{E}_{out} = \dot{q}_{con}$$

$$T_s = 92^{\circ}\text{C} + 0.3 \times 10^6 \frac{\text{W}}{\text{m}^3} \times 0.1\text{m} / 500\text{W}/\text{m}^2 \cdot \text{K} = 92^{\circ}\text{C} + 60^{\circ}\text{C} = 152^{\circ}\text{C}.$$

It follows that

$$T_o = 0.3 \times 10^6 \text{W}/\text{m}^3 \times (0.1\text{m})^2 / 2 \times 25\text{W}/\text{m} \cdot \text{K} + 152^{\circ}\text{C}$$

$$T_o = 60^{\circ}\text{C} + 152^{\circ}\text{C} = 212^{\circ}\text{C}.$$

## Exercise

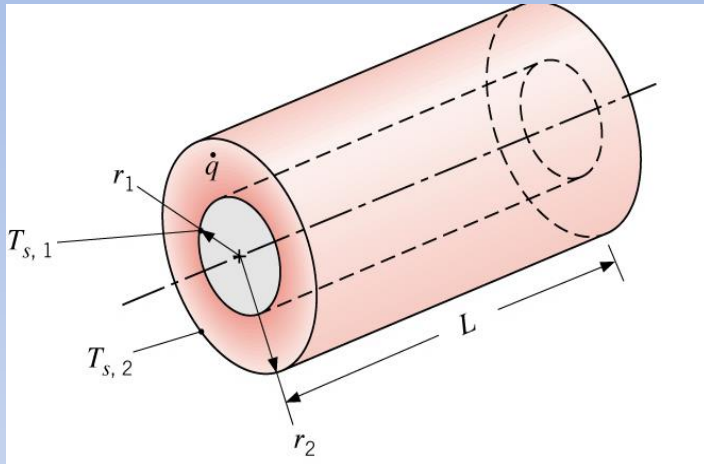
How can you find the heat flux leaving the wall?

Method 1

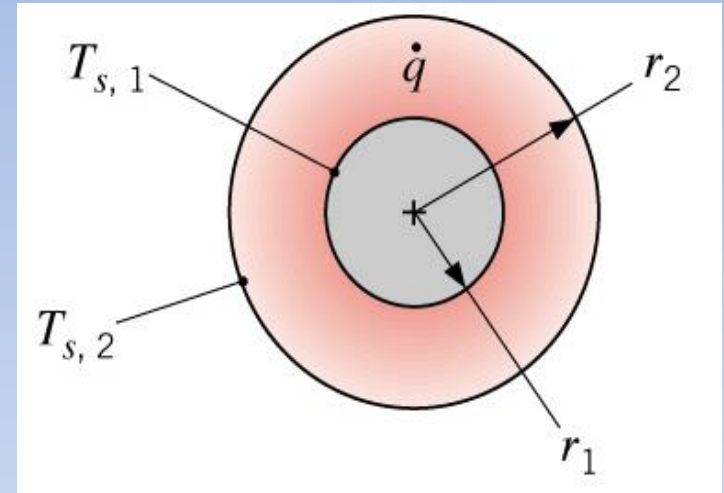
Method 2

# Radial Systems

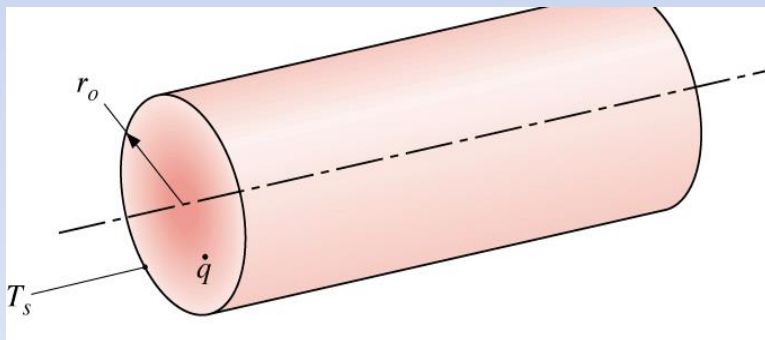
## Cylindrical (Tube) Wall



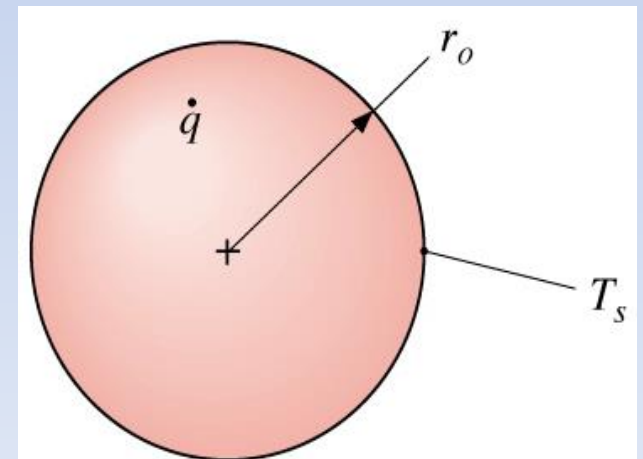
## Spherical Wall (Shell)



## Solid Cylinder (Circular Rod)

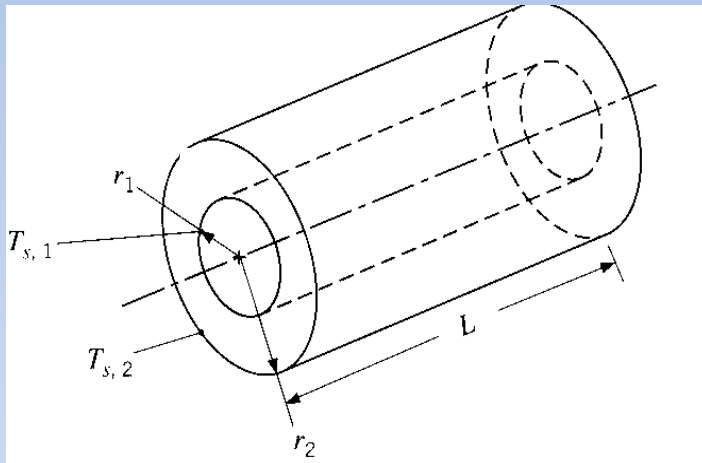


## Solid Sphere



# Hollow cylinder no generation only r-direction

## Cylindrical (Tube) Wall



Find the temp. distribution and the heat rate.

Follow the same procedure.

$$\frac{1}{r} \frac{\partial}{\partial r} \left( kr \frac{\partial T}{\partial r} \right) + \cancel{\frac{1}{r^2} \frac{\partial}{\partial \phi} \left( k \frac{\partial T}{\partial \phi} \right)} + \cancel{\frac{\partial}{\partial z} \left( k \frac{\partial T}{\partial z} \right)} + \cancel{\dot{q}} = \cancel{\rho c_p} \frac{\partial T}{\partial t}$$

$$\frac{1}{r} \frac{d}{dr} \left( kr \frac{dT}{dr} \right) = 0$$

# To obtain temp. distribution

- Solve the previous eq. and applying suitable B.Cs.
- Assuming the value of  $k$  to be constant. Double integration gives the general solution:

$$T(r) = C_1 \ln r + C_2$$

$$T(r_1) = T_{s,1} \quad \text{and} \quad T(r_2) = T_{s,2}$$

Applying these conditions to the general solution, we then obtain

$$T_{s,1} = C_1 \ln r_1 + C_2 \quad \text{and} \quad T_{s,2} = C_1 \ln r_2 + C_2$$

Solving for  $C_1$  and  $C_2$  and substituting into the general solution, we then obtain



$$T(r) = \frac{T_{s,1} - T_{s,2}}{\ln(r_1/r_2)} \ln\left(\frac{r}{r_2}\right) + T_{s,2}$$

Apply Fourier's Law

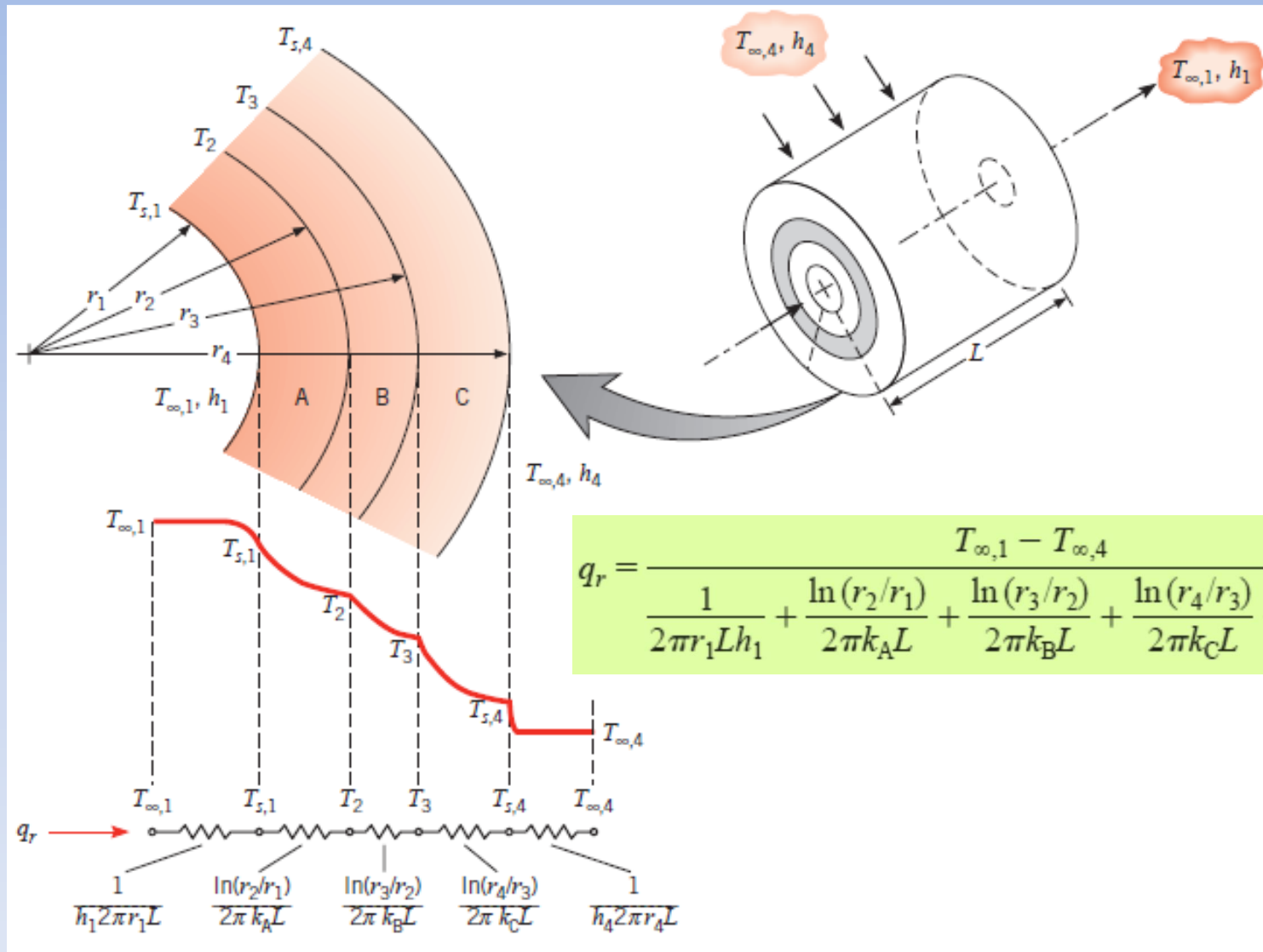
$$q_r = -kA \frac{dT}{dr} = -k(2\pi rL) \frac{dT}{dr}$$

$$q_r = \frac{2\pi Lk(T_{s,1} - T_{s,2})}{\ln(r_2/r_1)}$$

The thermal resistance is

$$R_{t,\text{cond}} = \frac{\ln(r_2/r_1)}{2\pi Lk}$$

# Composite cylindrical wall



# Note

$$q_r = \frac{T_{\infty,1} - T_{\infty,4}}{\frac{1}{2\pi r_1 L h_1} + \frac{\ln(r_2/r_1)}{2\pi k_A L} + \frac{\ln(r_3/r_2)}{2\pi k_B L} + \frac{\ln(r_4/r_3)}{2\pi k_C L} + \frac{1}{2\pi r_4 L h_4}} \quad (a)$$

Overall heat transfer coefficient,  $U$ , is obtained as follows

$$q_r = \frac{T_{\infty,1} - T_{\infty,4}}{R_{\text{tot}}} = UA(T_{\infty,1} - T_{\infty,4}) \quad (b)$$

If  $U$  is defined in terms of the inside area,  $A_1 = 2\pi r_1 L$ , the above eqs. (a) and (b) yield

$$U_1 = \frac{1}{\frac{1}{h_1} + \frac{r_1}{k_A} \ln \frac{r_2}{r_1} + \frac{r_1}{k_B} \ln \frac{r_3}{r_2} + \frac{r_1}{k_C} \ln \frac{r_4}{r_3} + \frac{r_1}{r_4} \frac{1}{h_4}}$$

# Look!!

This definition is *arbitrary*, and the overall coefficient may also be defined in terms of the outside surface area,  $A_4$  or any of the intermediate areas. Note that

$$U_1A_1 = U_2A_2 = U_3A_3 = U_4A_4 = (\Sigma R_t)^{-1}$$

and the specific forms of  $U_2$ ,  $U_3$ , and  $U_4$  may be inferred from Equations (a) and (b).

# Radial Systems

- Heat diffusion equation in the r-direction for steady-state conditions:

$$\frac{1}{r} \frac{d}{dr} \left( kr \frac{dT}{dr} \right) + \dot{q} = 0$$

- Assume  $k = \text{constant}$
- General Solution:

$$T = -\frac{\dot{q}}{4k} r^2 + C_1 \ln r + C_2$$

- Boundary Conditions:  $\left. \frac{dT}{dr} \right|_{r=0} = 0, \quad T(r_o) = T_s$

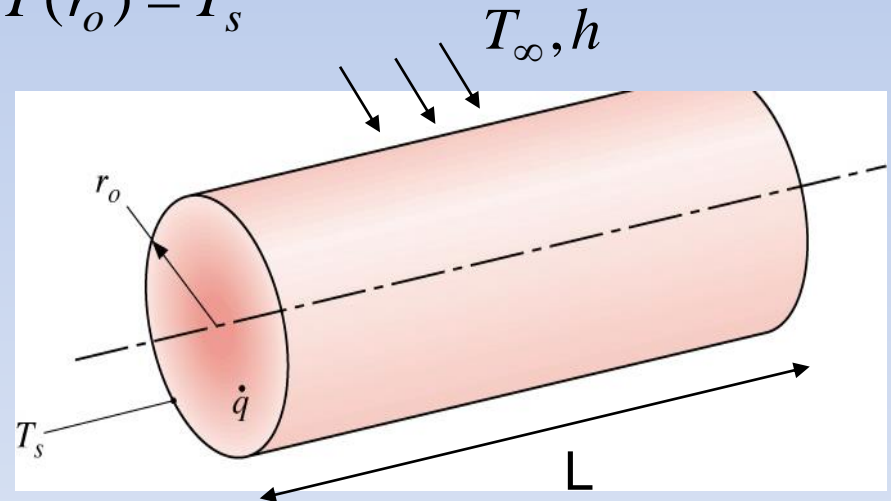
- Temperature profile:

$$T(r) = \frac{\dot{q} r_o^2}{4k} \left( 1 - \frac{r^2}{r_o^2} \right) + T_s \quad (3.6)$$

- Calculation of surface temperature:

$$\dot{q}(\pi r_o^2 L) = h(2\pi r_o L)(T_s - T_\infty) \quad \text{and}$$

$$T_s = T_\infty + \frac{\dot{q} r_o}{2h} \quad (3.7)$$



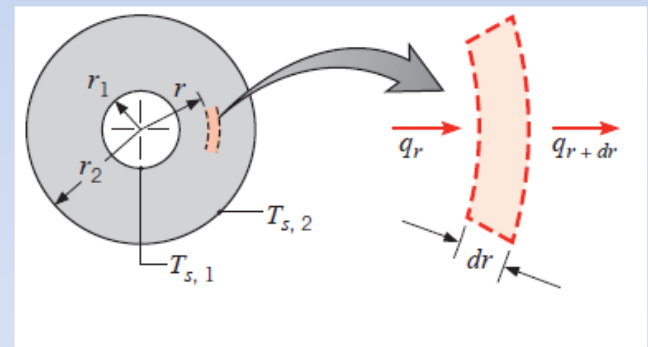
# For Spherical shape

- Fourier's law, for steady-state, one-dimensional conditions with no heat generation

$$q_r = -kA \frac{dT}{dr} = -k(4\pi r^2) \frac{dT}{dr}$$

- where  $A = 4\pi r^2$  is the area normal to the direction of heat transfer

$$\frac{q_r}{4\pi} \int_{r_1}^{r_2} \frac{dr}{r^2} = - \int_{T_{s,1}}^{T_{s,2}} k(T) dT$$





- Assuming  $k = \text{cons.}$

$$q_r = \frac{4\pi k(T_{s,1} - T_{s,2})}{(1/r_1) - (1/r_2)}$$

- Thermal resistance

$$R_{t,\text{cond}} = \frac{1}{4\pi k} \left( \frac{1}{r_1} - \frac{1}{r_2} \right)$$

- For temp. distribution follow the common procedure as explained before.

# Summary

## One-dimensional, steady-state solutions to the heat equation with no generation

	Plane Wall	Cylindrical Wall <sup>a</sup>	Spherical Wall <sup>a</sup>
Heat equation	$\frac{d^2T}{dx^2} = 0$	$\frac{1}{r} \frac{d}{dr} \left( r \frac{dT}{dr} \right) = 0$	$\frac{1}{r^2} \frac{d}{dr} \left( r^2 \frac{dT}{dr} \right) = 0$
Temperature distribution	$T_{s,1} - \Delta T \frac{x}{L}$	$T_{s,2} + \Delta T \frac{\ln(r/r_2)}{\ln(r_1/r_2)}$	$T_{s,1} - \Delta T \left[ \frac{1 - (r_1/r)}{1 - (r_1/r_2)} \right]$
Heat flux ( $q''$ )	$k \frac{\Delta T}{L}$	$\frac{k \Delta T}{r \ln(r_2/r_1)}$	$\frac{k \Delta T}{r^2 [(1/r_1) - (1/r_2)]}$
Heat rate ( $q$ )	$kA \frac{\Delta T}{L}$	$\frac{2\pi Lk \Delta T}{\ln(r_2/r_1)}$	$\frac{4\pi k \Delta T}{(1/r_1) - (1/r_2)}$
Thermal resistance ( $R_{t,cond}$ )	$\frac{L}{kA}$	$\frac{\ln(r_2/r_1)}{2\pi Lk}$	$\frac{(1/r_1) - (1/r_2)}{4\pi k}$